

Variational limits for active particles models

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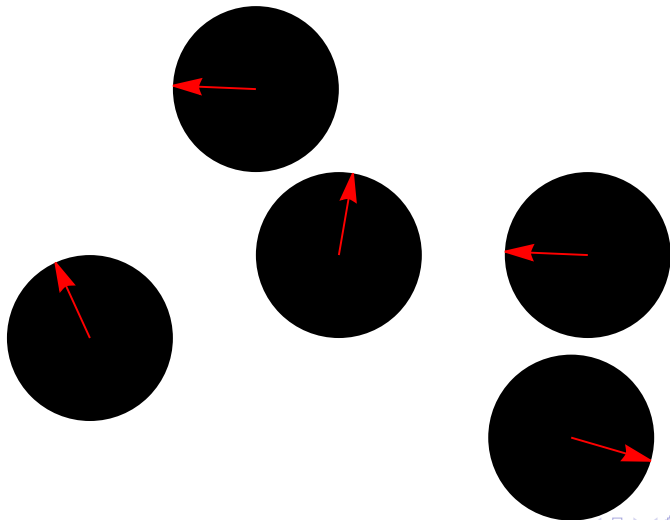
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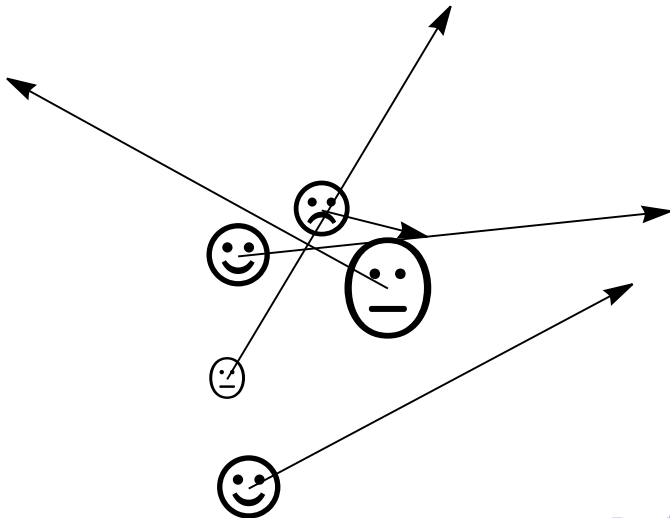


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Typical dynamics used to model active interacting systems is

$$\begin{cases} \dot{X}_i = v(X_i, \theta_i) \\ \dot{\theta}_i = F(X_i, \theta_i; \frac{1}{N} \sum_j \delta_{(X_j, \theta_j)}) + \dot{W}_i \end{cases}$$

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X will represent the space of spacial coordinates and Θ will denote the space of attached features.

A reference example is given by $X = \mathbb{R}^d$ and $\Theta = S^{d-1}$ (picture with black disks, sometimes called the Ising model of active particles):

$$\left\{ \begin{array}{l} \dot{X}_i = \hat{n}_{\theta_i} \\ \dot{\theta}_i = \frac{\sum_{j=1}^N \chi_r(X_i - X_j) f(\theta_i - \theta_j)}{\sum_{j=1}^N \chi_r(X_i - X_j)} + \dot{W}_i \end{array} \right.$$

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There are various informal predictions on the behavior of these kind of systems, as $N \rightarrow \infty$, the interaction radius r vanishes and a suitable long-time asymptotic is considered. See Degond et al.

It is a classical fact that, as $N \rightarrow \infty$, the empirical measure $\frac{1}{N} \sum_i \delta_{X_i, \theta_i}$ is *very* close to the solution to

$$\partial_t \mu_t = -d_x^*(\mu_t v) + (\Delta \mu_t + \operatorname{div}(\mu_t F[\mu_t])) + \sqrt{\mu_t/N} \operatorname{div}(\dot{W})$$

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or otherwise said, the large deviations rate functional of the law of the empirical measure is given by the control problem

$$\partial_t \mu_t + d_x^*(\mu_t v) - (\Delta \mu_t + \operatorname{div}(\mu_t F[\mu_t])) = \operatorname{div}(\mu_t E)$$

$$\operatorname{Cost}(\mu) = \frac{1}{2} \int dt \int d\mu_t(x) E^2(t, x)$$

Dynamic functional

Hereafter we assume that X and Θ are smooth, connected, compact manifolds without boundary. Θ is equipped with a Riemann metric and $d\theta$ is the associated normalized volume measure; but we stress that X does not have any natural metric, and it is only equipped with a reference probability measure M .

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Let $\mathcal{U} := \mathcal{C}([0, T]; \mathcal{P}(X \times \Theta))$. For our aims, the statistical properties of these systems are described by a dynamical free energy functional $I: \mathcal{U} \rightarrow [0, \infty]$ (see below)

$$I_1(\mu) = \left\| \partial_t \mu_t + d_x^*(\mu_t \mathbf{v}) - (\Delta \mu_t + \operatorname{div}(\mu_t F[\mu_t])) \right\|_{L_2([0, T]; H_\mu^{-1})}^2$$

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where

- The interaction $F = F[\mu](x, \theta)$ writes as

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$$I_1(\mu) = \left\| \partial_t \mu_t + d_x^*(\mu_t v) - (\Delta \mu_t + \operatorname{div}(\mu_t F[\mu_t])) \right\|_{L_2([0, T]; H_\mu^{-1})}^2$$

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- The equilibrium measure is $dm(x, \theta) = e^{-\psi(x, \theta)} dM(x) d\theta$.
- $v = v(x, \theta)$ is a vector field on X , such that $d_x^* v = 0$ and $m(v) = 0$.
- The initial data is fixed.

As suggested in the literature, we consider the scaling

$$I_\varepsilon(\mu) = \left\| \varepsilon \partial_t \mu_t + d_x^*(\mu_t v) - \varepsilon^{-1} (\Delta \mu_t + \operatorname{div}(\mu_t F[\mu_t])) \right\|_{L_2([0, T]; H_\mu^{-1})}^2$$

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The functional I_ε can be understood as the optimal control problem

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Most importantly, one is interested in a limit in a variational sense, which is relevant for LD and control.

Understanding the scaling

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For each fixed x , the second equation admits an invariant measure μ_x on Θ . As $\varepsilon \rightarrow 0$, X converges to the solution to $\dot{X} = \bar{V}(X)$, where

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But if $\bar{V} = 0$, rather consider the limit of $X_{t/\varepsilon}$:

$$\begin{aligned}\dot{X} &= \frac{1}{\varepsilon} v(X, \theta) \\ \dot{\theta} &= \frac{1}{\varepsilon^2} F(X, \theta) + \frac{1}{\varepsilon} \dot{W}\end{aligned}$$

We want to identify the Γ -limit of I_ε as $\varepsilon \rightarrow 0$. In the optimal control interpretation, this means that we want to calculate the limiting cost of a $\mu \in \mathcal{U}$ as the minimal limiting cost

$$I(\mu) = \inf_{\mu_\varepsilon \rightarrow \mu} \lim_{\varepsilon} \int dt d\mu_{\varepsilon,t} |E^\varepsilon|^2$$

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From the LD point of view, this means that as $N \rightarrow \infty$ and then $\varepsilon \rightarrow 0$

$$\mathbb{P}(\pi_{N,\varepsilon} \sim \mu) = \exp(-N I(\mu))$$

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I_ε is equicoercive and it holds $\Gamma\text{-}\lim_\varepsilon I_\varepsilon = I$, where I is defined in the next slide.

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This implies that minima and minimizers converge, that constrained problems converge, that the inequalities hold in the limit of LD etc.

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$I: \mathcal{U} \rightarrow [0, \infty]$ is defined as follows, for suitable tensors A and B .

$$I(\mu) := \begin{cases} \|\partial_t \varrho - d_x^*(B[\varrho]d_x \varrho)\|_{L_2([0, T]; H_{A[\varrho]}^{-1})}^2 & \text{if } \mu_t = \varrho_t m \\ +\infty & \text{otherwise} \end{cases}$$

This has again a control interpretation, but in the x -variable only. If

$$\partial_t \varrho - d_x^*(B[\varrho]d_x \varrho) = d_x^*(A[\varrho]E)$$

then $I(\varrho m) = \int dt dm A[\varrho] E \cdot E$.

Einstein's relations

For each smooth 1-form ω on X , there is a unique solution φ such that $m(\varphi) = 0$ and

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For ϱ a probability density on X , the tensors A and B are given by

$$A[\varrho](\omega, \omega') := m(\varrho \nabla \langle \Phi[\varrho], \omega \rangle \cdot \nabla \langle \Phi[\varrho], \omega' \rangle)$$

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The Einstein's relation $B[\varrho] = h''(\varrho)A[\varrho]$ holds, with $h(\varrho) = \varrho \log \varrho$.

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- For instance, on \mathbb{R}^d , with constant v and *non-interacting particles*, one gathers $A(\varrho) = \varrho$ and this reduces to the JKO formulation on the heat equation.
- If $d_x^* v \neq 0$, a drift term appears.

Hint of the proof

Consider the rescaled control problem

$$\varepsilon \partial_t \mu_t^\varepsilon + \mathbf{d}_x^*(\mu_t^\varepsilon \mathbf{v}) = \varepsilon^{-1} \Delta \mu_t^\varepsilon + \varepsilon^{-1} \operatorname{div}(\mu_t^\varepsilon \mathbf{F}[\mu_t^\varepsilon]) + \operatorname{div}(\mu^\varepsilon \mathbf{E}^\varepsilon)$$

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This implies that the marginal in x satisfies a continuity equation

$$\partial_t \varrho_t^\varepsilon + \varepsilon^{-1} \int_{\Theta} d\mu_t^\varepsilon \mathbf{v} = 0$$

and the current $\varepsilon^{-1} \int_{\Theta} d\mu_t^\varepsilon \mathbf{v}$ is order 1 as $\varepsilon \downarrow 0$. However $d_x^*(\mu_t^\varepsilon \mathbf{v})$ is not order ε .

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- taking θ -average $\varepsilon \partial_t \mu_t^\varepsilon \rightarrow 0$.
- $d_x^*(\mu_t^\varepsilon \mathbf{v}) \rightarrow (\partial_t \varrho) m$.
- When tested again $\varphi = \langle \Phi[\varrho], \omega \rangle$, the terms order ε^{-1} converge to the controlled forcing of the limiting functional.

Grazie!

